

From (quantum) geometry to spectral theory

Juan Margalef Bentabol

Joint work with F. Barbero and Eduardo J.S. Villaseñor



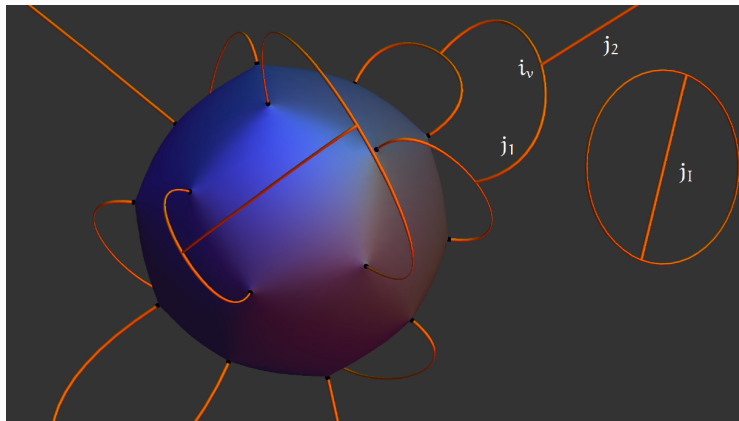
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Introduction

Goal

To understand some features of the spectrum of the area operator



Eigenvalues of the (simplified) area operator

$$a_s = \sum_{n_i \in M_N} \sqrt{n_i(n_i + 2)} \quad \left| \begin{array}{l} M_N := \{n_1, \dots, n_N\} \in (\mathbb{N}^+)^N \\ N \in \mathbb{N} \end{array} \right.$$

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$$M_N = \{\underbrace{1, \dots, 1}_{N_1}, \underbrace{2, \dots, 2}_{N_2}, \underbrace{3, \dots, 3}_{N_3}, \dots\} \equiv (N_1, N_2, \dots) \in c_{00}(\mathbb{N})$$

Eigenvalues of the (simplified) area operator

eigenvalue	$(N_1, N_2, \dots)$
0	$(0, 0, 0, 0 \dots)$
$\sqrt{3}$	$(1, 0, 0, 0 \dots)$
$2\sqrt{2}$	$(0, 1, 0, 0 \dots)$
$2\sqrt{3}$	$(2, 0, 0, 0 \dots)$
$\sqrt{15}$	$(0, 0, 1, 0 \dots)$
$2\sqrt{2} + \sqrt{3}$	$(1, 1, 0, 0 \dots)$
$2\sqrt{6}$	$(0, 0, 0, 1 \dots)$
$3\sqrt{3}$	$(3, 0, 0, 0 \dots)$
$\sqrt{3} + \sqrt{15}$	$(1, 0, 1, 0 \dots)$

eigenvalue	$(N_1, N_2, \dots)$
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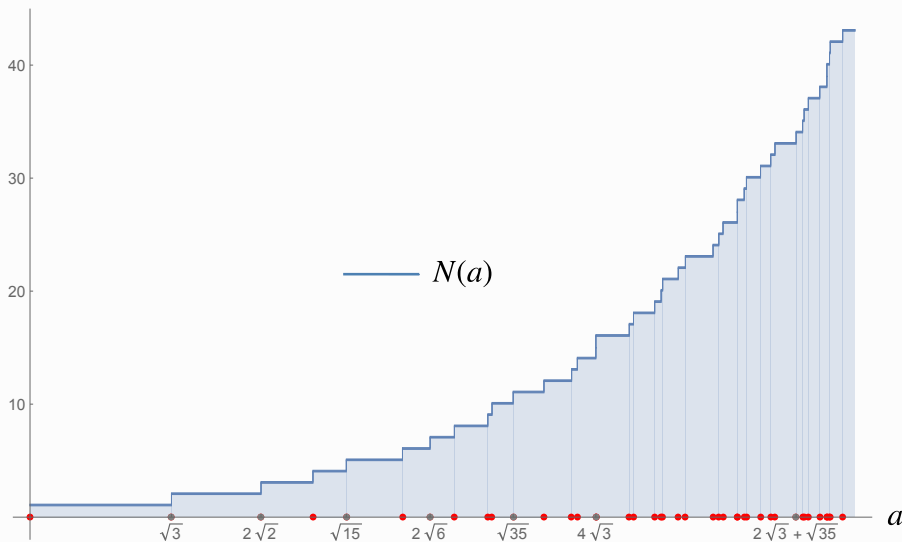
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$$N(a) := 1 + \# \left\{ M_N \in c_{00}(\mathbb{N}) \quad / \quad \sum_{n_i \in M_N} \sqrt{n_i(n_i + 2)} \leq a \right\}$$

Eigenvalues of the (simplified) area operator



First approach: approximations

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$$\sqrt{x(x+2)} \sim x+1 - \frac{1}{2x} + O\left(\frac{1}{x^2}\right) \quad \longrightarrow \quad \sqrt{n_i(n_i+2)} \sim n_i+1$$

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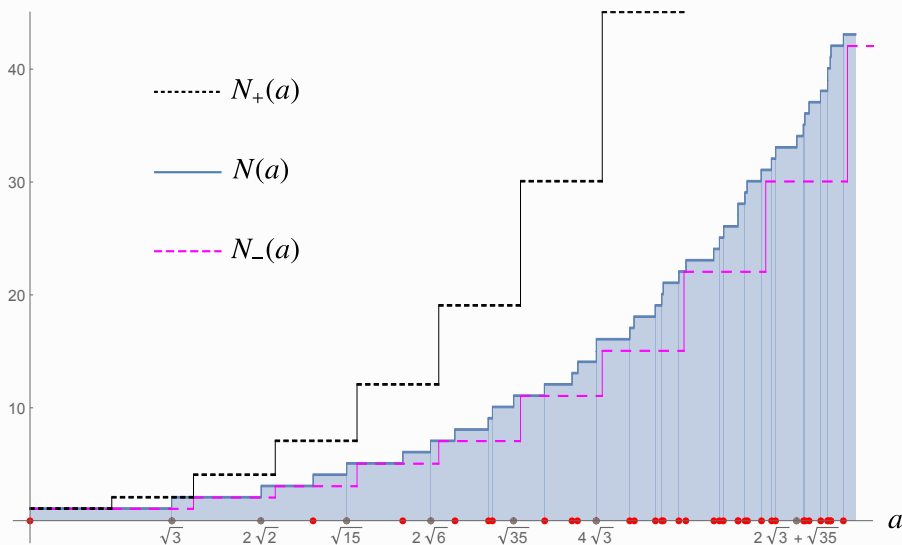
$$\sum n_i < \sum \sqrt{n_i(n_i+2)} < \sum (n_i+1) \\ N_-(a) \leq N(a) \leq N_+(a)$$

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Playing with partitions

Using generating functions and ideas of Hardy, Ramanujan, Littlewood, Rademacher. . . we obtain

$$N_-(a) \sim \frac{1}{4a\sqrt{3}} \text{Exp} \left(\pi \sqrt{\frac{2a}{3}} \right) \quad N_+(a) \sim \frac{1}{2\pi\sqrt{2a}} \text{Exp} \left(\pi \sqrt{\frac{2a}{3}} \right)$$

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- Same exponent factor but different prefactor.
- We can use $n_i + \varepsilon < \sqrt{n_i(n_i + 2)} < n_i + 1$ for $\varepsilon \in (0, 1)$.
- No clear approximation for N . New ideas are needed!

Second approach: Laplace transform

N as a staircase function

We rewrite $N(a) = \sum_{n=1}^{\infty} \theta(a - a_n)$ $a_n = \sum s_k \in \text{Spec}(N)$

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$$\begin{aligned} N(a) &= \frac{1}{2\pi i} \int_{c-i\infty}^{c+i\infty} \sum_{n=1}^{\infty} \frac{e^{-z(a-a_n)}}{z} dz = \frac{1}{2\pi i} \int_{c-i\infty}^{c+i\infty} \frac{e^{za}}{z} \sum_{n=1}^{\infty} e^{-za_n} dz = \dots \\ &= \frac{1}{2\pi i} \int_{c-i\infty}^{c+i\infty} \frac{e^{za}}{z} \text{Exp} \left[- \sum_{k=1}^{\infty} \ln(1 - e^{-zs_k}) \right] dz - \theta(a) \end{aligned}$$

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- $s_k = \sqrt{k(k+2)}$ our problem at hand

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$$N(a) \sim \frac{1}{2\pi^{3/2}i} \int_{c-i\infty}^{c+i\infty} z^{1/2} \text{Exp} \left((a - a_*)z + \frac{\pi^2}{6z} + \frac{z}{4} \ln z \right) F(z) dz$$

Approximating N

$$N(a) \sim \frac{1}{\sqrt{\pi}} \alpha^{3/2} e^{-\alpha/2} F(\alpha) I_{-3/2} \left(\frac{\pi^2}{3\alpha} + \frac{\alpha}{4} \right) =: N_0(a)$$

$$\alpha \sim \frac{\pi}{\sqrt{6a}} \quad (a \rightarrow \infty) \quad I_{-3/2}(z) = \sqrt{\frac{2}{\pi}} \frac{z \sinh z - \cosh z}{z^{3/2}}$$

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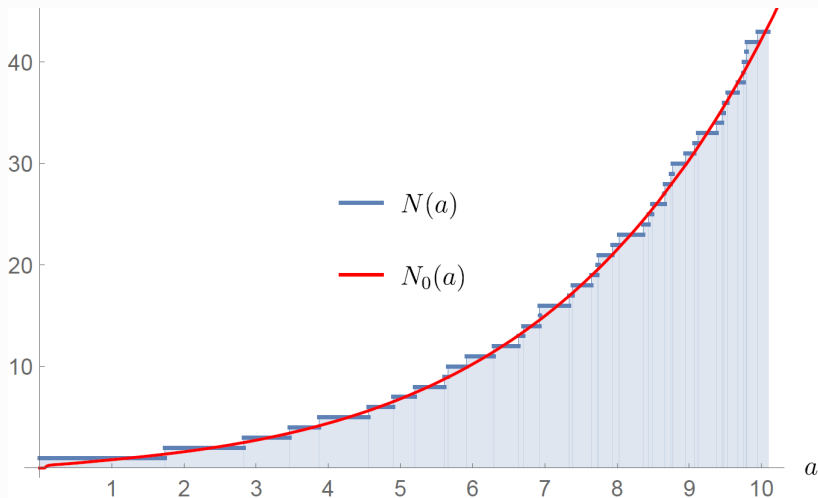
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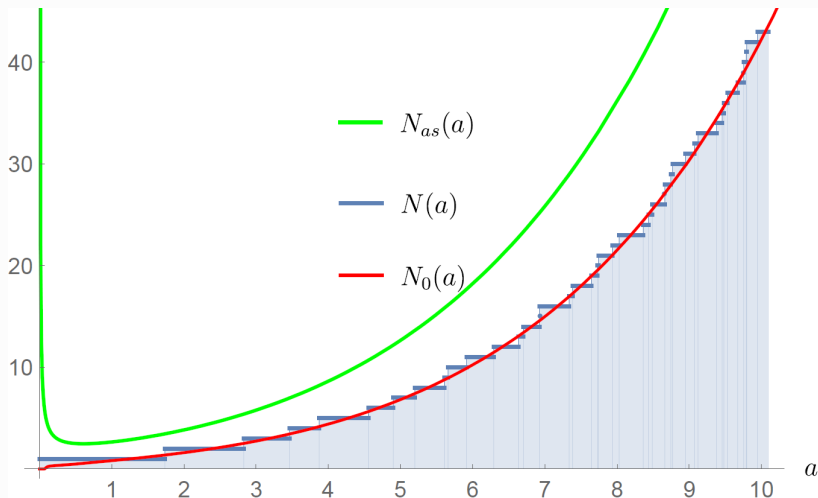
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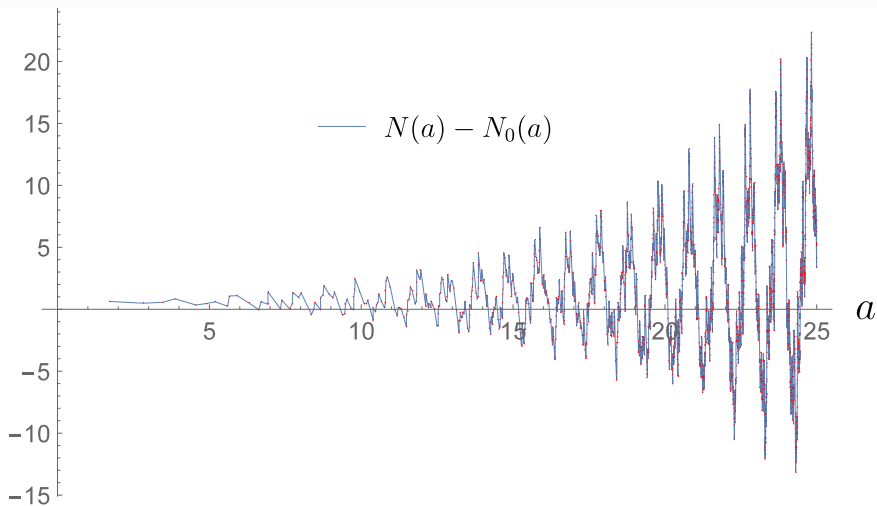
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- $N_-(a)$ is really good, it only fails by a **constant** $\sqrt{2}$ factor. Difficult to guess *a priori*.
- We get a very accurate representation for $N(a)$ valid for all areas.
- The used methods may be useful to tackle similar problems.

Thanks for your attention



J. Fernando Barbero, Juan Margalef-Bentabol, Eduardo J.S. Villaseñor, *On the distribution of the eigenvalues of the area operator in loop quantum gravity*. CQG 35 (2018), [arXiv:1712.06918]



M. Barreira, M. Carfora, C. Rovelli, *Physics with nonperturbative quantum gravity: Radiation from a quantum black hole*. GRG 28 (1996).



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I. Agullo, F. Barbero, E.F. Borja, J. Diaz-Polo, E.J.S. Villaseñor, *Detailed black hole state counting in loop quantum gravity*. PRD 82 (2010)



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N as a staircase function

a	$N_{as}(a)$	$N_0(a)$	$N(a)$
-----	-------------	----------	--------

10	68	42	43
30	8599	6282	6282
50	307719	237955	237788
70	6099037	4870521	4867770