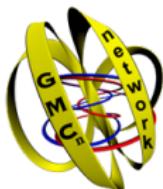


Short talk: Numerical optimal control of nonholonomic mechanical systems.

Rodrigo T. Sato Martín de Almagro

Supervisor:

David Martín de Diego



Dec, 2018 - 13YRWGMC



The author has been partially supported by Ministerio de Economía, Industria y Competitividad (MINEICO, Spain) under grants MTM 2015-64166-C2-2P, MTM 2016-76702-P.

Table of Contents

1 Introduction

2 Variational integrators

- High-order variational integrators
- Variational integrators and optimal control

3 Nonholonomic integrators and optimal control

- Nonholonomic mechanics
- Nonholonomic optimal control

Table of Contents

1 Introduction

2 Variational integrators

- High-order variational integrators
- Variational integrators and optimal control

3 Nonholonomic integrators and optimal control

- Nonholonomic mechanics
- Nonholonomic optimal control

Mechanical systems

Mechanical systems

- Described via Lagrangian or Hamiltonian formulation
- Built-in geometric properties (Manifold structure of configuration space, symplecticity)
- Built-in conservation laws due to symmetries (Noether theorem)

Symplectic integrators

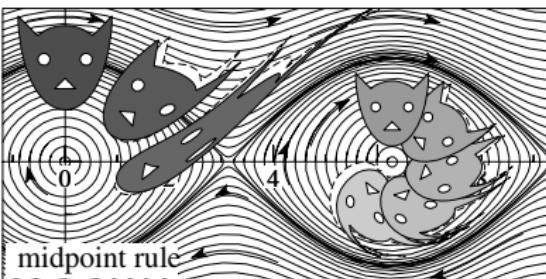
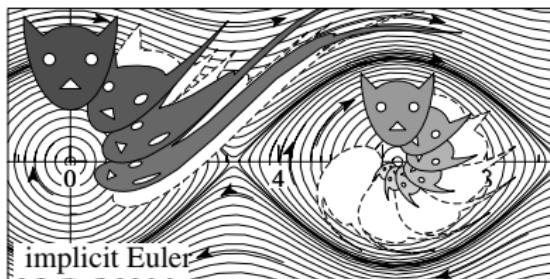
Why do we like symplectic integrators?

- Good qualitative and quantitative behaviour.

Symplectic integrators

Why do we like symplectic integrators?

- Good qualitative and quantitative behaviour.
- Preserve flow properties (symplecticity, momenta...).



Symplectic integrators

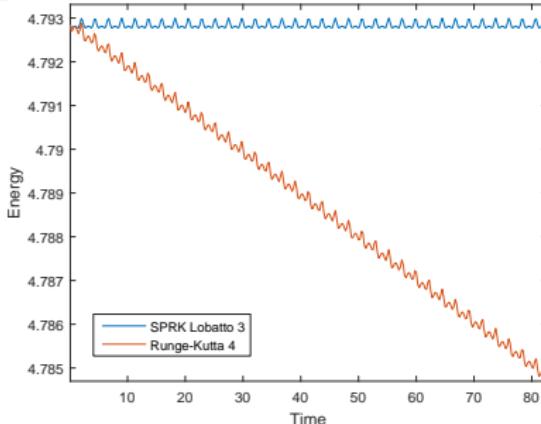
Why do we like symplectic integrators?

- Good qualitative and quantitative behaviour.
- Preserve flow properties (symplecticity, momenta...).
- Energy not exactly preserved... [Ge & Marsden]

Symplectic integrators

Why do we like symplectic integrators?

- Good qualitative and quantitative behaviour.
- Preserve flow properties (symplecticity, momenta...).
- Energy not exactly preserved... [Ge & Marsden]
- ... but good long-term energy behaviour.



Symplectic integrators

Why do we like symplectic integrators?

- Good qualitative and quantitative behaviour.
- Preserve flow properties (symplecticity, momenta...).
- Energy not exactly preserved... [Ge & Marsden]
- ... but good long-term energy behaviour.

How come energy behaves so well?

Theorems ([Moser], [Benettin & Giorgilli], [Tang], [Murua]...) warrant that symplectic integrators are integrating exactly some existing Hamiltonian system that is *close* to the original one.

Table of Contents

1 Introduction

2 Variational integrators

- High-order variational integrators
- Variational integrators and optimal control

3 Nonholonomic integrators and optimal control

- Nonholonomic mechanics
- Nonholonomic optimal control

Generating symplectic integrators easily

Variational integrators are always symplectic.

Idea ([Veselov], [Suris], [Marsden & West]...)

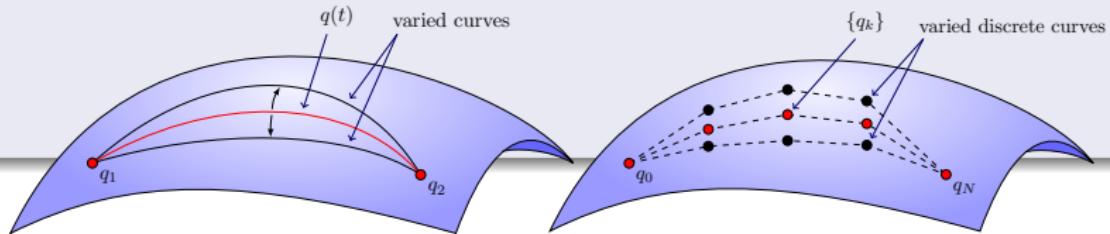
- Substitute continuous state space with discrete one.

Generating symplectic integrators easily

Variational integrators are always symplectic.

Idea ([Veselov], [Suris], [Marsden & West]...)

- Substitute continuous state space with discrete one.



Generating symplectic integrators easily

Variational integrators are always symplectic.

Idea ([Veselov], [Suris], [Marsden & West]...)

- Substitute continuous state space with discrete one.
- Build discrete analogue of Hamilton's principle.

Generating symplectic integrators easily

Variational integrators are always symplectic.

Idea ([Veselov], [Suris], [Marsden & West]...)

- Substitute continuous state space with discrete one.
- Build discrete analogue of Hamilton's principle.
- Derive equations of motion and conservations from the principle.

Generating symplectic integrators easily

Variational integrators are always symplectic.

Idea ([Veselov], [Suris], [Marsden & West]...)

- Substitute continuous state space with discrete one.
- Build discrete analogue of Hamilton's principle.
- Derive equations of motion and conservations from the principle.

Discrete equations of motion = Difference equations (a.k.a. **our integrator**).

The starting point

Hamilton-Pontryagin action

$(q, v, p) : [0, T] \subset \mathbb{R} \rightarrow TQ \oplus T^*Q$, sufficiently differentiable curve and fixed boundary values $q(0) = q_a$, $q(T) = q_b$.

$$\mathcal{J}_{\mathcal{HP}}(q, v, p) = \int_0^T [L(q(t), v(t)) + \langle p(t), \dot{q}(t) - v(t) \rangle] dt$$

Dynamical equations

$$\frac{dq(t)}{dt} = v(t),$$

$$\frac{dp(t)}{dt} = D_1 L(q(t), v(t)),$$

$$p(t) = D_2 L(q(t), v(t))$$

Variationally partitioned Runge-Kutta integrators

Discrete Hamilton-Pontryagin action

$$(\mathcal{J}_{\mathcal{HP}})_d = \sum_{k=0}^{N-1} \sum_{i=1}^s h b_i \left[L(Q_k^i, V_k^i) + \left\langle P_k^i, \frac{Q_k^i - q_k}{h} - \sum_{j=1}^s a_{ij} V_k^j \right\rangle \right. \\ \left. + \left\langle p_{k+1}, \frac{q_{k+1} - q_k}{h} - \sum_{j=1}^s b_j V_k^j \right\rangle \right]$$

where (a_{ij}, b_j) coefficients of a Runge-Kutta (RK) method.

Discrete dynamics

Discrete dynamical equations: Symplectic partitioned RK methods

$$q_{k+1} = q_k + h \sum_{j=1}^s b_j V_k^j, \quad D_2 L(q_{k+1}, v_{k+1}) = D_2 L(q_k, v_k) + h \sum_{i=1}^s \hat{b}_j D_1 L(Q_k^j, V_k^j),$$

$$Q_k^i = q_k + h \sum_{j=1}^s a_{ij} V_k^j, \quad D_2 L(Q_k^i, V_k^i) = D_2 L(q_k, v_k) + h \sum_{j=1}^s \hat{a}_{ij} D_1 L(Q_k^j, V_k^j),$$

where $(\hat{a}_{ij}, \hat{b}_j)$ satisfy $b_i \hat{a}_{ij} + \hat{b}_j a_{ji} = b_i \hat{b}_j$ and $\hat{b}_i = b_i$.

Here we have used the continuous relation $p(t) = D_2 L(q(t), v(t))$ to relate p_k and p_{k+1} with v_k and v_{k+1} .

Controlled mechanical systems

Controlled dynamics

Let $f : TQ \times U \rightarrow T^*Q$ define a mechanical forcing, $U \subset \mathbb{R}^k$ control space.

$$\frac{dq(t)}{dt} = v(t),$$

$$\frac{dp(t)}{dt} = D_1 L(q(t), v(t)) + f(q(t), v(t), u(t)),$$

$$p(t) = D_2 L(q(t), v(t))$$

$(q, v, u) : [0, T] \subset \mathbb{R} \rightarrow TQ \times U$. Controlled eqs. of motion define section of $TT^*Q \times U$, of the form $(q, G(q, v), v, F(q, v, u), u)$.

Action principle for optimal control

Cost function

$$C : TQ \times U \rightarrow \mathbb{R}$$

Hamilton-Pontryagin action

$(q, p, \xi_q, \xi_p, \mu_q, \mu_p, u) : [0, T] \subset \mathbb{R} \rightarrow TT^*Q \oplus T^*T^*Q \times U$
 sufficiently differentiable curve with fixed boundary values
 $(q(0), p(0)) = (q_a, p_a), (q(T), p(T)) = (q_b, p_b).$

$$\begin{aligned} \mathcal{J}_{HP}(q, p, \xi_q, \xi_p, \mu_q, \mu_p, u) = & \int_0^T [C(q(t), \xi_q(t), u(t)) \\ & + \langle (\mu_q(t), \mu_p(t)), (\dot{q}(t) - \xi_q(t), \dot{p}(t) - \xi_p(t)) \rangle] dt \end{aligned}$$

evaluated over the mechanical section $(q, G(q, v), v, F(q, v, u), \mu_q, \mu_p, u).$

Optimal control

Necessary conditions for optimality (Dynamical equations)

$$\dot{\mu}_q + \dot{\mu}_p D_1 G(q, v) = D_1 C(q, v, u) - \mu_p D_1 F(q, v, u),$$

$$\dot{\mu}_p D_2 G(q, v) = D_2 C(q, v, u) - \mu_p D_2 F(q, v, u) - \mu_q,$$

$$0 = D_3 C(q, v, u) - \mu_p D_3 F(q, v, u),$$

$$\dot{q} = v,$$

$$D_1 G(q, v) \dot{q} + D_2 G(q, v) \dot{v} = F(q, v, u).$$

Discrete optimal control

Discrete Hamilton-Pontryagin action

$$\begin{aligned}
 \mathcal{J}_d = & \sum_{k=0}^{N-1} \sum_{i=1}^s h b_i \left[C \left(Q_k^i, V_k^i, U_k^i \right) \right. \\
 & + \left\langle M_{Q,k}^i, \frac{Q_k^i - q_k}{h} - \sum_{j=1}^s a_{ij} V_k^j \right\rangle + \left\langle \mu_{q,k+1}, \frac{q_{k+1} - q_k}{h} - \sum_{j=1}^s b_j V_k^j \right\rangle \\
 & + \left\langle M_{P,k}^i, \frac{D_2 L(Q_k^i, V_k^i) - D_2 L(q_k, v_k)}{h} - \sum_{j=1}^s \hat{a}_{ij} \left[D_1 L(Q_k^j, V_k^j) + f(Q_k^j, V_k^j, U_k^j) \right] \right\rangle \\
 & \left. + \left\langle \mu_{p,k+1}, \frac{D_2 L(q_{k+1}, v_{k+1}) - D_2 L(q_k, v_k)}{h} - \sum_{j=1}^s \hat{b}_j \left[D_1 L(Q_k^j, V_k^j) + f(Q_k^j, V_k^j, U_k^j) \right] \right\rangle \right]
 \end{aligned}$$

Taking variations we find discrete necessary conditions for optimality.

Table of Contents

1 Introduction

2 Variational integrators

- High-order variational integrators
- Variational integrators and optimal control

3 Nonholonomic integrators and optimal control

- Nonholonomic mechanics
- Nonholonomic optimal control

The setting

Nonholonomic Lagrangian systems

(L, Q, N) , where $N \subset TQ$ constraint manifold. Locally described by null-set of $\Phi : TQ \rightarrow \mathbb{R}^m$, $m = \text{codim}_{TQ} N$.

Dynamical equations

Obtained via Chetaev's principle, not variational.

$$\begin{aligned}\dot{q} &= v, \\ \dot{p} &= D_1 L(q, v) + \langle \lambda, D_2 \Phi(q, v) \rangle, \\ p &= D_2 L(q, v), \\ 0 &= \Phi(q, v).\end{aligned}$$

where λ Lagrange multipliers.

The setting (continued)

Nonholonomic controlled Lagrangian systems

(L, Q, N, f, U) , where $N \subset TQ$ constraint manifold. Locally described by null-set of $\Phi : TQ \rightarrow \mathbb{R}^m$, $m = \text{codim}_{TQ} N$.

Dynamical equations

Obtained via Chetaev's + D'Alembert's principle.

$$\dot{q} = v,$$

$$\dot{p} = D_1 L(q, v) + \langle \lambda, D_2 \Phi(q, v) \rangle + f(q, v, u),$$

$$p = D_2 L(q, v),$$

$$0 = \Phi(q, v).$$

where λ Lagrange multipliers.

From constrained to free and forced

Lagrange multipliers

If the system is regular, then it is possible to differentiate $\Phi = 0$ w.r.t. t , substitute equations and solve for λ , leading to

$$\lambda(t) = I(q, v), \quad \text{or} \quad \lambda(t) = I(q, v, u)$$

It ensures every initial condition on N remains on N . This transforms the constrained system into a free forced system.

Controlled nonholonomic eqs. of motion still define section of $TT^*Q \times U$, of the form $(q, G(q, v), v, F(q, v, u), u)$ with $F(q, v, u) = D_1L(q, v) + \langle I(q, v, u), D_2\Phi(q, v) \rangle + f(q, v, u)$.

Nonholonomic integrator

For Lobatto-type RK methods with (a_{ij}, b_j) :

Nonholonomic integrator (finally on arxiv:1810.10926)

$$\begin{aligned}
 q_{k+1} &= q_k + h \sum_{i=1}^s b_i V_k^i, \quad D_2 L(q_{k+1}, v_{k+1}) = D_2 L(q_k, v_k) + h \sum_{i=1}^s \hat{b}_i \left[D_1 L(Q_k^i, V_k^i) + \left\langle \Lambda_k^i, D_2 \Phi(Q_k^i, V_k^i) \right\rangle \right], \\
 Q_k^i &= q_k + h \sum_{j=1}^s a_{ij} V_k^j, \quad D_2 L(Q_k^i, V_k^i) = D_2 L(q_k, v_k) + h \sum_{j=1}^s \hat{a}_{ij} \left[D_1 L(Q_k^i, V_k^j) + \left\langle \Lambda_k^j, D_2 \Phi(Q_k^i, V_k^j) \right\rangle \right], \\
 q_k^i &= Q_k^i, \quad D_2 L(q_k^i, v_k^i) = D_2 L(q_k, v_k) + h \sum_{j=1}^s a_{ij} \left[D_1 L(Q_k^i, V_k^j) + \left\langle \Lambda_k^j, D_2 \Phi(Q_k^i, V_k^j) \right\rangle \right], \\
 0 &= \Phi(q_k^i, v_k^i).
 \end{aligned}$$

Hamiltonian flow:

$F_{L_d}^\Lambda : TQ|_N \times \Lambda \rightarrow TQ|_N \times \Lambda$, $(q_k, v_k, \lambda_k) \mapsto (q_{k+1}, v_{k+1}, \lambda_{k+1})$,
 with $\lambda_k = \Lambda_k^1, \lambda_{k+1} = \Lambda_k^s$.

Lagrange multipliers in the discrete setting

Here comes trouble!

- If $\Lambda_k^i \mapsto I(Q_k^i, V_k^i, U_k^i)$, constraint not preserved in general!

Lagrange multipliers in the discrete setting

Here comes trouble!

- If $\Lambda_k^i \mapsto I(Q_k^i, V_k^i, U_k^i)$, constraint not preserved in general!
- Cannot easily solve for $\{\Lambda_k^i\}_{i=2}^s$, contrary to continuous case!

Lagrange multipliers in the discrete setting

Here comes trouble!

- If $\Lambda_k^i \mapsto I(Q_k^i, V_k^i, U_k^i)$, constraint not preserved in general!
- Cannot easily solve for $\{\Lambda_k^i\}_{i=2}^s$, contrary to continuous case!
- If solved, $\{\Lambda_k^i\}_{i=2}^s$ must be a function of $\lambda_k = \Lambda_k^1$ too!

Lagrange multipliers in the discrete setting

Here comes trouble!

- If $\Lambda_k^i \mapsto I(Q_k^i, V_k^i, U_k^i)$, constraint not preserved in general!
- Cannot easily solve for $\{\Lambda_k^i\}_{i=2}^s$, contrary to continuous case!
- If solved, $\{\Lambda_k^i\}_{i=2}^s$ must be a function of $\lambda_k = \Lambda_k^1$ too!

Solution

Imposing variations on $\{\Lambda_k^i\}_{i=2}^s$ appropriately.

Can also be done in the continuous case, but *a priori* might seem overkill.

Treating multipliers as independent variables

Continuous Hamilton-Pontryagin action

$$\begin{aligned}\mathcal{J}_{\mathcal{HP}} = \int_0^T & \left[C(q(t), v(t), u(t)) \right. \\ & + \langle (\mu_q, \mu_p), (\dot{q} - v, (G(q, v))^\cdot - F(q, v, u, \lambda)) \rangle \\ & \left. + \langle \mu_\lambda, \lambda - I(q, v, u) \rangle \right] dt\end{aligned}$$

Treating multipliers as independent variables (continued)

Necessary conditions for optimality (Dynamical equations)

$$\dot{\mu}_q + \dot{\mu}_p D_1 G(q, v) = D_1 C(q, v, u) - \mu_p D_1 F(q, v, u, \lambda) - \mu_\lambda D_1 I(q, v, u),$$

$$\dot{\mu}_p D_2 G(q, v) = D_2 C(q, v, u) - \mu_p D_2 F(q, v, u, \lambda) - \mu_\lambda D_2 I(q, v, u) - \mu_q,$$

$$0 = D_3 C(q, v, u) - \mu_p D_3 F(q, v, u, \lambda) - \mu_\lambda D_3 I(q, v, u),$$

$$0 = \mu_\lambda - \mu_p D_4 F(q, v, u, \lambda),$$

$$\dot{q} = v,$$

$$D_1 G(q, v) \dot{q} + D_2 G(q, v) \dot{v} = F(q, v, u, \lambda),$$

$$\lambda = I(q, v, u).$$

Treating multipliers as independent variables (continued)

Necessary conditions for optimality (Dynamical equations)

$$\dot{\mu}_q + \dot{\mu}_p D_1 G(q, v) = D_1 C(q, v, u) - \mu_p D_1 F(q, v, u, \lambda) - \mu_\lambda D_1 I(q, v, u),$$

$$\dot{\mu}_p D_2 G(q, v) = D_2 C(q, v, u) - \mu_p D_2 F(q, v, u, \lambda) - \mu_\lambda D_2 I(q, v, u) - \mu_q,$$

$$0 = D_3 C(q, v, u) - \mu_p D_3 F(q, v, u, \lambda) - \mu_\lambda D_3 I(q, v, u),$$

$$0 = \mu_\lambda - \mu_p D_4 F(q, v, u, \lambda),$$

$$\dot{q} = v,$$

$$D_1 G(q, v) \dot{q} + D_2 G(q, v) \dot{v} = F(q, v, u, \lambda),$$

$$0 = \Phi(q, v).$$

Only need derivatives of $I(q, v, u)$. Equivalent to imposing
 $\delta \lambda = D_1 I(q, v, u) \delta q + D_2 I(q, v, u) \delta v + D_3 I(q, v, u) \delta u$.

Lagrange multipliers in the discrete setting

Discrete Lagrange multiplier variations

Casting nonholonomic integrator in the form

$$\mathcal{F}(x, y) = 0,$$

where $x = (q_k, v_k, \lambda_k, U_k^1, \dots, U_k^s)$, $y = (q_{k+1}, v_{k+1}, \lambda_{k+1}, \Lambda_k^2, \dots, \Lambda_k^{s-1}, \dots)$, then

$$\delta y = - (D_2 \mathcal{F}(x, y))^{-1} D_1 \mathcal{F}(x, y) \delta x$$

We get $\{\delta \Lambda_k^i\}_{i=2}^s$ in terms of $\delta q_k, \delta v_k, \delta \lambda_k, \{\delta U_k^i\}_{i=1}^s$.

Discrete nonholonomic optimal control

Discrete Hamilton-Pontryagin action

$$\mathcal{J}_d =$$

$$\begin{aligned} & \sum_{k=0}^{N-1} \sum_{i=1}^s h b_i \left[C \left(Q_k^i, V_k^i, U_k^i \right) + \left\langle M_{Q,k}^i, \frac{Q_k^i - q_k}{h} - \sum_{j=1}^s a_{ij} V_k^j \right\rangle + \left\langle \mu_{q,k+1}, \frac{q_{k+1} - q_k}{h} - \sum_{j=1}^s b_j V_k^j \right\rangle \right. \\ & + \left\langle M_{P,k}^i, \frac{D_2 L(Q_k^i, V_k^i) - D_2 L(q_k, v_k)}{h} - \sum_{j=1}^s \hat{a}_{ij} \left[D_1 L(Q_k^j, V_k^j) + f(Q_k^j, V_k^j, U_k^j) + \Lambda_k^j D_2 \Phi(Q_k^j, V_k^j) \right] \right\rangle \\ & \left. + \left\langle \mu_{p,k+1}, \frac{D_2 L(q_{k+1}, v_{k+1}) - D_2 L(q_k, v_k)}{h} - \sum_{j=1}^s \hat{b}_j \left[D_1 L(Q_k^j, V_k^j) + f(Q_k^j, V_k^j, U_k^j) + \Lambda_k^j D_2 \Phi(Q_k^j, V_k^j) \right] \right\rangle \right] \end{aligned}$$

Imposition of variations of Lagrange multipliers

$$\langle \delta \mathcal{J}_d, \delta c_d \rangle + \sum_{k=0}^{N-1} \sum_{i=2}^s h b_i \left\langle M_{\Lambda,k}^i, \delta \Lambda_k^i - R_k^{q,i} \delta q_k - R_k^{v,i} \delta v_k - R_k^i \delta \lambda_k - \sum_{j=1}^s R_k^{U^j,i} \delta U_k^j \right\rangle$$

Discrete nonholonomic optimal control II

Closing the system

Resulting necessary conditions system must be supplemented with

$$D_2 L(q_k^i, v_k^i) = D_2 L(q_k, v_k)$$

$$+ h \sum_{j=1}^s a_{ij} \left[D_1 L(Q_k^j, V_k^j) + f(Q_k^j, V_k^j, U_k^j) + \mathcal{N}_k^j D_2 \Phi(Q_k^j, V_k^j) \right],$$

$$q_k^i = Q_k^i,$$

$$0 = \Phi(q_k^i, v_k^i).$$

Example: Nonholonomic controlled particle

$$Q = \mathbb{R}^3$$

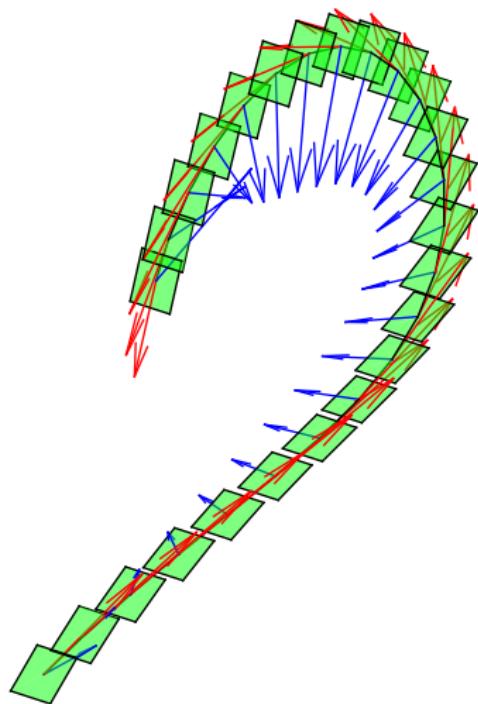
$$U = \mathbb{R}^2$$

$$L(q, v) = \frac{1}{2} (v_x^2 + v_y^2 + v_z^2)$$

$$\Phi(q, v) = v_z - yv_x$$

$$f(q, v, u) = u^1 dx + u^2 dy + yu^1 dz$$

$$C(q, v, u) = \frac{1}{2} (u_1^2 + u_2^2) .$$



THANKS FOR YOUR
ATTENTION!